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## DIFFERENTIAL EQUATIONS WITH INFINITE NUMBER OF LIMIT CYCLES

Consider a system of differential equations

$$(1) \quad \dot{x}_1 = f_1, \quad \dot{x}_2 = f_2$$

with  $f_i$  sufficiently regular, say of class  $C^2$ , functions defined on the (phase-)plane  $R^2$ . Let  $x_i(t)$ ,  $i = 1, 2$ , denote a periodic solution of (1), and  $K$  — the corresponding orbit which is a closed curve in  $R^2$ .

$K$  is called a *limit cycle* if no closed orbit other than  $K$  lies in its sufficiently small neighbourhood. If all orbits near  $K$  wind toward  $K$  for  $t \rightarrow +\infty$  (resp.,  $t \rightarrow -\infty$ ), the limit cycle  $K$  is called *stable* (resp., *unstable*). In the case where orbits lying on one side of  $K$  wind toward  $K$  for  $t \rightarrow +\infty$  and those lying on the other side wind toward  $K$  for  $t \rightarrow -\infty$ ,  $K$  is called a *semi-stable* limit cycle. For instance, the well-known Van der Pol equation

$$\dot{x} = y, \quad \dot{y} = (1 - x^2)y - x$$

has a stable limit cycle which is seen in the phase diagram (see Fig. 1) obtained by modelling this equation on an analogue computer.

Recently, D'Heedene [1] presented a full analysis of the equation of the form

$$(2) \quad \ddot{x} + \mu \sin \dot{x} + x = 0.$$

It was shown that (2) has limit cycles in infinite number. Moreover, for each of them it was possible to provide a characterization of stability type.

The careful examination shows that the methods and results in [1] may be generalized to the case of the class of equations

$$(3) \quad \ddot{x} + \mu h(\dot{x}) + x = 0$$

with  $h$  satisfying the following conditions:

- $$(4) \quad \begin{array}{ll} \text{(i)} & h \in C^1(R), \\ \text{(ii)} & \text{periodic (with period } 2T), \\ \text{(iii)} & h(-s) = -h(s), \\ \text{(iv)} & h(s) > 0 \text{ for } 0 < s < T. \end{array}$$

Equation (3) is equivalent to the system

$$\dot{x} = y, \quad \dot{y} = -x - \mu h(y),$$

the analysis of which, due to the symmetry of the system, may be reduced to the half-plane  $y \geq 0$ .

Denote by  $M$  a positive constant such that  $|h| \leq M$  and  $|\dot{h}| \leq M$ . Then the following theorem is true:

**THEOREM I.** *For any  $\mu$ ,  $|\mu| < 2/M$ , equation (3) with  $h$  satisfying (4) has infinitely many limit cycles displaced on the plane according to the following rule: for any  $n = 1, 2, \dots$  there is exactly one limit cycle among*

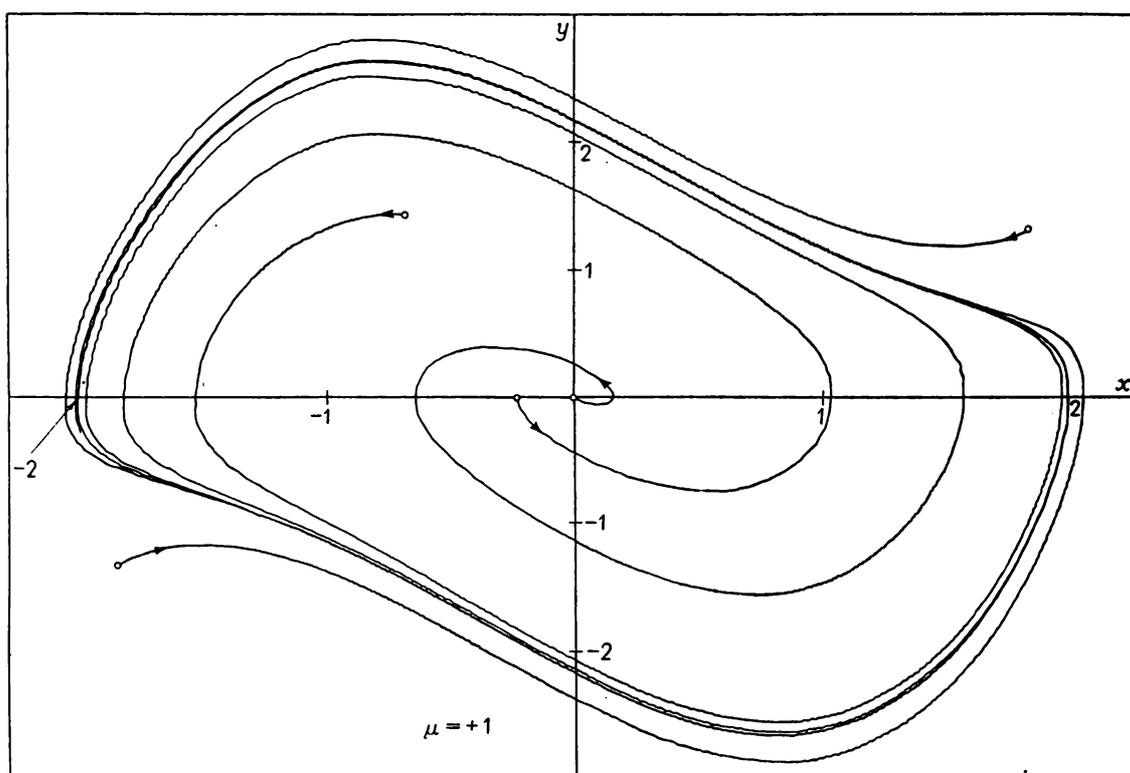


Fig. 1

all orbits passing through the points  $(\mu h(y), y)$ ,  $y \in (nT, (n+1)T)$ , of the plane. Moreover, for  $\mu > 0$  the limit cycles which correspond to  $n$  even are stable while those corresponding to  $n$  odd are unstable. For  $\mu < 0$  the opposite holds.

Condition  $|\mu| < 2/M$  in Theorem I may be relaxed.

**THEOREM II.** *Conclusions of Theorem I hold true for those  $n$  and  $\mu$  which satisfy the inequality*

$$n > \frac{2}{T} M |\mu| [M |\mu| + 1 + (M^2 \mu^2 + 2M |\mu|)^{1/2}].$$

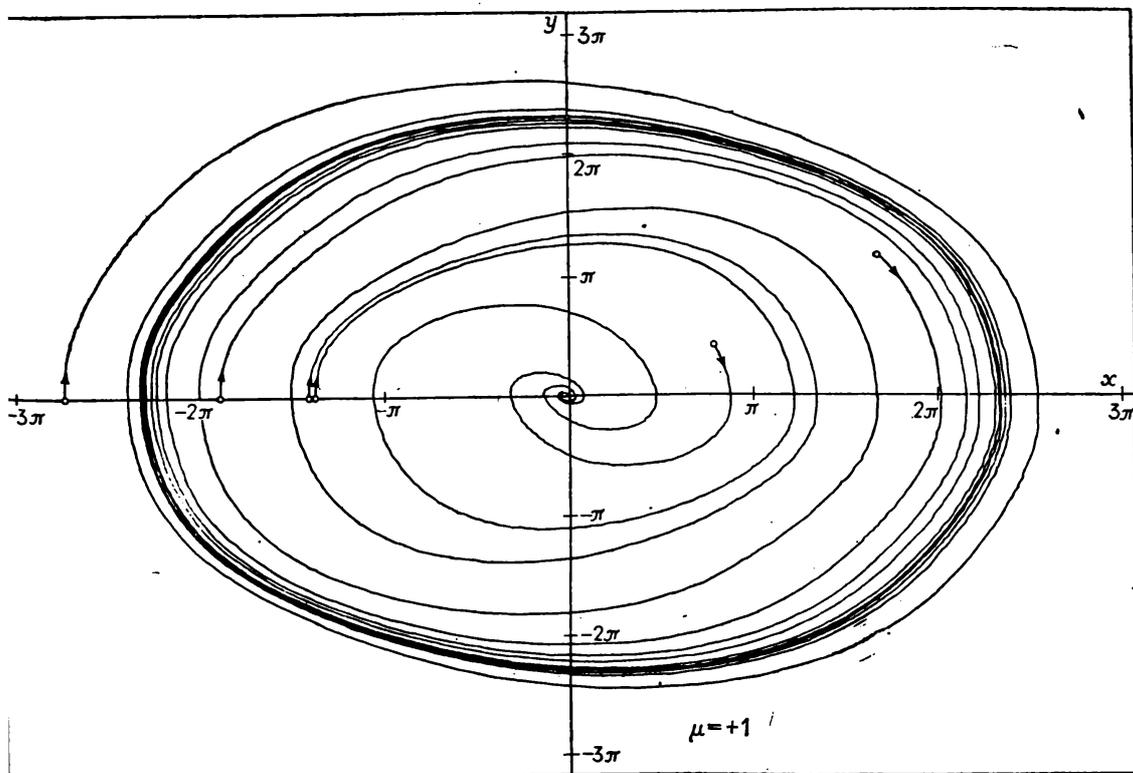
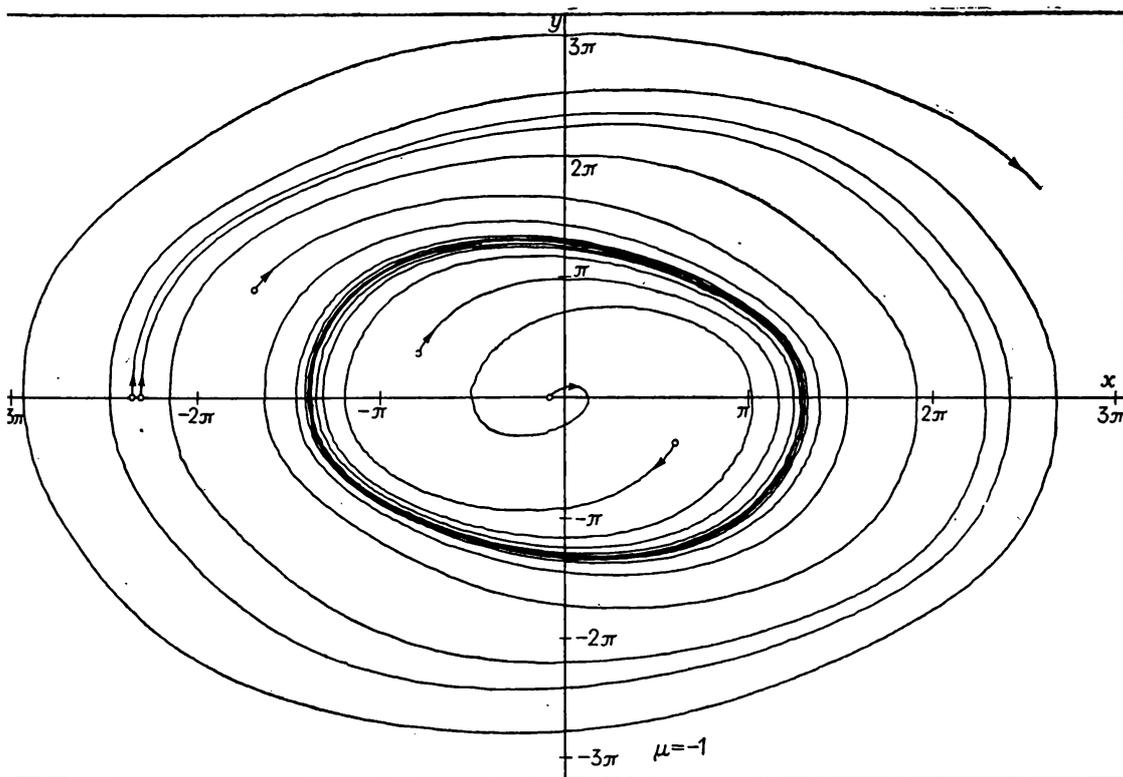


Fig. 2

The proofs of Theorems I and II are parallel to those presented in [1] with minor modifications only. Therefore we omit them.

It is clear that equation (2) is a special case of (3). There we have  $T = \pi$ , and  $M = 1$ . In Fig. 2 there are given two diagrams for equation (2) with  $\mu = +1$  and  $\mu = -1$ . The diagrams were obtained by modelling the equation on an analogue computer (like those in Fig. 1). For technical reasons the range of variable  $y$  was restricted to the interval  $(0, 3\pi)$ . We can observe that, in agreement with the theorems, the following holds:

- (a) no limit cycle passes through the segment  $x = 0$ ,  $0 \leq y \leq \pi$ ;
- (b) through the segment  $x = 0$ ,  $\pi \leq y \leq 2\pi$  there passes exactly one limit cycle which is unstable for  $\mu = 1$  and stable for  $\mu = -1$ ;
- (c) through the segment  $x = 0$ ,  $2\pi \leq y \leq 3\pi$  there passes exactly one limit cycle which is stable for  $\mu = 1$  and unstable for  $\mu = -1$ .

Initial points for orbits in our diagrams were arbitrarily chosen. The behaviour of orbits exhibits clearly the existence of a limit cycle: in the case of a stable (unstable) limit cycle the orbits wind toward (resp., outward) it.

There is no method as far to answer the question concerning the existence of limit cycles in a way which would be practically useful for all equations. Thus the numerical computations and the application of analogue computers seem to be the only general means for analyzing such problems.

#### Reference

- [1] R. M. D'Heedene, *For all real  $\mu$ ,  $\ddot{x} + \mu \sin \dot{x} + x = 0$  has an infinite number of limit cycles*, Journal of Differential Equations 5 (1969), p. 564-571.

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### RÓWNANIA RÓŻNICZKOWE Z NIESKOŃCZONĄ LICZBĄ CYKLI GRANICZNYCH

#### STRESZCZENIE

Praca przedstawia uogólnienie wyniku D'Heedene'a, mówiącego o istnieniu nieskończenie wielu cykli granicznych dla równania (2) na ogólniejszą klasę równań

(3). W pracy nie podano dowodu tego faktu, gdyż jest on jedynie rachunkową modyfikacją oryginalnego dowodu D'Heedene'a.

Ponadto przedstawiono wykresy otrzymane przez zaprogramowanie na maszynie analogowej równania (2) rozpatrywanego przez D'Heedene'a i równania Van der Pola. Jest to w nie których przypadkach jedyna metoda pozwalająca rozstrzygnąć istnienie cyklu granicznego.

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